

Chemistry 430/530a Practice

Midterm 1

Ziad Ganim 9/21/2015

Problem 1.1.

$$k_B T = \frac{200 \text{ cm}^{-1}, 0.6 \text{ kcal/mol}, 2.5 \text{ kJ/mol}}{\text{(any unit system)}}$$

$$Q(N, V, T) = \frac{\sum_j e^{-\beta E_j}}{\quad}$$

$$Q(N, V, E) = \frac{1}{N}$$

$$S(p_j) = \frac{-k_B \sum_j p_j \ln p_j}{\quad}$$

$$\binom{N}{n} = \frac{N!}{(N-n)! n!}$$

(how many ways to order N distinguishable particles in n slots)

Problem 1.2.

In class we derived, $\langle E \rangle$ and $\langle E^2 \rangle$. Please derive an expression for the ensemble average energy to the n th power, $\langle E^n \rangle$ using your $Q(N, V, T)$ from Problem 1.1.

$$\text{Since } \langle E \rangle = -\frac{1}{Q} \frac{\partial Q}{\partial \beta} \text{ and } \langle E^2 \rangle = \frac{1}{Q} \frac{\partial^2 Q}{\partial \beta^2}$$

$$\text{Try } \langle E^n \rangle \stackrel{?}{=} \frac{\partial^{(n)} Q}{\partial \beta^{(n)}}$$

$$\frac{\partial^{(n)} Q}{\partial \beta^{(n)}} = \sum_j \frac{\partial^{(n)}}{\partial \beta^{(n)}} e^{-\beta E_j}$$

$$= \sum_j (-E_j) \frac{\partial^{(n-1)}}{\partial \beta^{(n-1)}} e^{-\beta E_j}$$

$$= \sum_j (-E_j)^2 \frac{\partial^{(n-2)}}{\partial \beta^{(n-2)}} e^{-\beta E_j}$$

$$\vdots$$

$$= \sum_j (-E_j)^{n-1} \frac{\partial}{\partial \beta} e^{-\beta E_j}$$

$$= \sum_j (-E_j)^n e^{-\beta E_j} = Q (-1)^n \langle E^n \rangle$$

$$\therefore \langle E^n \rangle = \frac{(-1)^n}{Q} \frac{\partial^{(n)} Q}{\partial \beta^{(n)}}$$

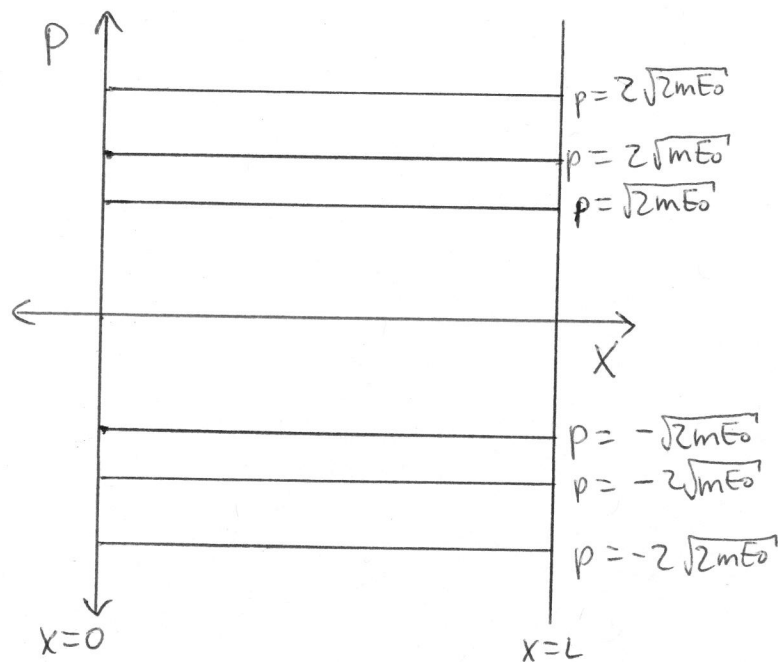
Problem 1.3.

Please draw, roughly to scale, the phase space plot for a classical particle in a 1D box of length L with energies E_0 , $2E_0$ and $4E_0$. Please label the axes. How does the distance between contours depend on energy?

Particle in a 1D Box: $X = [0, L]$

$$\frac{p^2}{2m} = E$$

$$p = \sqrt{2mE}$$



$$\frac{p_1}{p_2} = \sqrt{\frac{2mE_1}{2mE_2}} = \sqrt{\frac{E_1}{E_2}}$$

Problem 1.4.

Given

$$Q = \frac{e^{-\beta \hbar \omega / 2}}{1 - e^{-\beta \hbar \omega}},$$

please evaluate

$$A = -\frac{1}{\beta} \ln Q$$

$$\langle E \rangle = -\frac{\partial}{\partial \beta} \ln Q$$

$$S = k \ln Q + \frac{\langle E \rangle}{T}$$

$$C_V = \frac{\partial \langle E \rangle}{\partial T},$$

and show that $C_V = k_B$ when $k_B T \gg \hbar \omega$ using the approximation valid at high temperatures, $e^{\hbar \omega / k_B T} \approx 1 + \frac{\hbar \omega}{k_B T}$.

$$\ln Q = -\frac{\beta \hbar \omega}{2} - \ln[1 - e^{-\beta \hbar \omega}]$$

$$A = \frac{\hbar \omega}{2} + \frac{1}{\beta} \ln[1 - e^{-\beta \hbar \omega}]$$

$$\begin{aligned} \langle E \rangle &= -\frac{\hbar \omega}{2} + [1 - e^{-\beta \hbar \omega}]^{-1} \frac{\partial}{\partial \beta} [1 - e^{-\beta \hbar \omega}] \\ &= \frac{\hbar \omega}{2} + \frac{\hbar \omega e^{-\beta \hbar \omega}}{1 - e^{-\beta \hbar \omega}} = \frac{\hbar \omega}{2} + \frac{\hbar \omega}{e^{\beta \hbar \omega} - 1} \end{aligned}$$

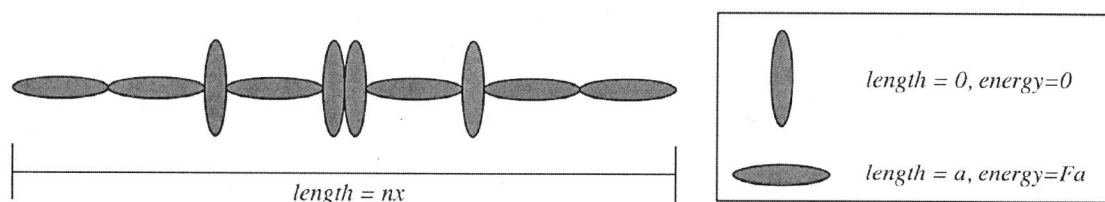
$$\begin{aligned} S &= -\frac{k \beta \hbar \omega}{2} - k \ln[1 - e^{-\beta \hbar \omega}] + \frac{\hbar \omega}{2T} + \frac{\hbar \omega}{T} \frac{1}{e^{\beta \hbar \omega} - 1} \\ &= -k \ln[1 - e^{-\beta \hbar \omega}] + \frac{\hbar \omega k}{\beta} \frac{1}{e^{\beta \hbar \omega} - 1} \end{aligned}$$

$$C_V = \frac{\partial \langle E \rangle}{\partial T} = \frac{\partial \langle E \rangle}{\partial \beta} \cdot \frac{\partial \beta}{\partial T} = \frac{\partial \langle E \rangle}{\partial \beta} \cdot \left(-\frac{1}{k T^2} \right)$$

$$= -\frac{\hbar \omega}{k T^2} [-(e^{\beta \hbar \omega} - 1)^{-2}] e^{\beta \hbar \omega} \cdot \hbar \omega = \frac{\hbar^2 \omega^2}{k T^2} \frac{e^{\beta \hbar \omega}}{(e^{\beta \hbar \omega} - 1)^2}$$

As $k T \gg \hbar \omega$ or $\beta \ll \hbar \omega$, $C_V \rightarrow \frac{\hbar^2 \omega^2}{k T^2} \frac{(1 + \beta \hbar \omega)}{(\beta \hbar \omega)^2} = \frac{\hbar^2 \omega^2}{k T^2 \beta^2 \hbar^2 \omega^2} + \frac{\hbar^2 \omega^2}{k T^2 \beta \hbar \omega} \rightarrow 0$
 $= k$

Problem 1.5. Here we consider a one-dimensional chain of n segments, as illustrated below. Each segment can exist in two states- length $x = a$ or length $x = 0$. When tension F is applied, the energy of a segment is $F \cdot x$, or $a \cdot F = aF$ for extended or $0 \cdot F = 0$ for contracted segments, respectively. Thus by changing F , we can change the energy difference between the two states.



What is the partition function, Q , for one segment at temperature T and tension F ?

$$Q = 1 + e^{-\beta Fa} = 1 + e^{-Fa/kT}$$

Using this result, calculate the average length, $\langle x \rangle$ of one segment as a function of T .

$$\langle x \rangle = \frac{(0)1 + (a)e^{-Fa/kT}}{1 + e^{-Fa/kT}} = \frac{a e^{-Fa/kT}}{1 + e^{-Fa/kT}}$$

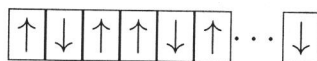
For n segments, the length will be $n\langle x \rangle$. Verify that this result leads to Hooke's Law, $n\langle x \rangle = c_0 + c_1 F$, at high temperatures where the limit $e^{-Fa/k_B T} \approx 1 - \frac{Fa}{k_B T}$ is applicable.

$$n\langle x \rangle = \frac{na e^{-Fa/kT}}{1 + e^{-Fa/kT}}$$

$$\text{As } kT \gg Fa, \quad n\langle x \rangle \rightarrow \frac{na(1 - Fa/kT)}{(2 - Fa/kT)^0} = \frac{na}{2} - \frac{na^2 F}{2kT}$$

$$c_0 = \frac{na}{2}, \quad c_1 = -\frac{na^2}{2kT} \quad (F < 0 \text{ for spring to stretch with increasing tension.})$$

Problem 1.6. Recall the 1D Ising model with N spins, each situated at a lattice site, and each spin may be up or down.



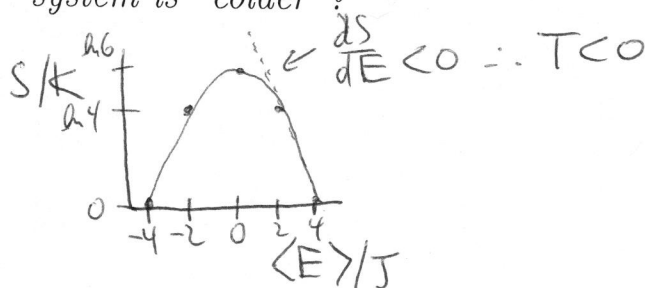
We will treat four spins with no nearest neighbor coupling. If we represent an \uparrow with $s_i = +1$ and \downarrow with $s_i = -1$, then the energy function for the system, the Hamiltonian, is

$$H = -J \sum_{i=1}^4 s_i. \quad (1)$$

For **four spins**, please write out all possible energy levels and the degeneracy for each energy level.

States	Energy	Ω	S/k
$\uparrow\uparrow\uparrow\uparrow$	$-4J$	1	$\ln 1 = 0$
$\uparrow\uparrow\uparrow\downarrow$ $\uparrow\uparrow\downarrow\uparrow$ $\uparrow\downarrow\uparrow\uparrow$ $\downarrow\uparrow\uparrow\uparrow$	$-2J$	4	$\ln 4$
$\uparrow\uparrow\downarrow\downarrow$ $\uparrow\downarrow\uparrow\downarrow$ $\downarrow\uparrow\downarrow\uparrow$ $\downarrow\downarrow\uparrow\uparrow$ $\downarrow\downarrow\downarrow\uparrow$ $\downarrow\uparrow\downarrow\downarrow$	0	6	$\ln 6$
$\downarrow\downarrow\downarrow\uparrow$ $\downarrow\downarrow\uparrow\downarrow$ $\downarrow\uparrow\downarrow\downarrow$ $\uparrow\downarrow\downarrow\downarrow$	$2J$	4	$\ln 4$
$\downarrow\downarrow\downarrow\downarrow$	$4J$	1	$\ln 1$

Sketch a plot of how the entropy S changes as a function of the energy of the system. Recall that we define the temperature $\frac{1}{T} = \frac{dS}{dE}$. Show that the temperature is negative as $E \rightarrow E_{max}$. Does negative temperature mean the system is "colder"?



These are non-thermal configurations with high internal energy — not "cold" in the conventional sense.